# Heat transfer study from square cylinder immersed in water-based nanofluid: Effect of orientation 

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#### Abstract

Heat transfer study from the heated square cylinder at a different orientation angle to the stream of nanofluids has been investigated numerically. CuO-based nanofluids were used to elucidate the significant effect of parameters: Reynolds number (1-40), nanoparticle volume fraction ( $0.00-0.05$ ), the diameter of the NPs $(30-100 \mathrm{mn})$ and the orientation of square cylinder $\left(0-90^{\circ}\right)$. The numerical results were expressed in terms of isotherm contours and average Nusselt number to explain the effect of relevant parameters. Over the range of conditions, the separation of the boundary layers of nanofluids increased with the size of the NPs as compared to pure water. NPs volume fraction and its size had a significant effect on heat transfer rate. The square cylinder of orientation angle $\left(45^{\circ}\right)$ gained a more efficient heat transfer cylinder than other orientation angles. Finally, the correlations were developed for the average Nusselt number in terms of the relevant parameters for $45^{\circ}$ orientation of the cylinder for new applications.


Keywords: inclined square cylinder, nanofluids, Reynolds number, Prandtl number, steady flow regime

## 1. INTRODUCTION

Nanofluids have gained a lot of attention from researchers due to many applications in various fields such as energy, bio, and pharmaceutical industry and chemical, environmental, materials, medical and thermal engineering. Nanofluids have gained significant attention due to their enhanced thermal characteristics as compared to base fluids. Over the last few decades, the heat transfer from a bluff body of different shapes to nanofluids has been frequently used in many industrial applications, for instance, pin-type and tubular heat exchangers, cooling of nuclear fuel rods and various electronic components, etc. Among various configurations of the bluff body, the circular cylinder has been studied more extensively due to its frequent occurrence in many practical applicationswhich is found in Chhabra and Richardson (2011). In contrast to the circular cylinder, not only knowledge of the square cross-section of the cylinder is very limited, but also there is additional complexity arising from an orientation of the cylinder with respect to the incoming

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nanofluids. All the reported results related to heat transfer indicated the varying levels of augmentation in heat transfer under appropriate conditions. For instance, Valipour and Ghadi (2011) have numerically investigated the flow and heat transfer over a circular cylinder for copper-based aqueous nanofluids in the range of $\operatorname{Re} \leq 40$ and $\phi \leq 0.05$. In the results part, the Nusselt number is seen to increase with the volume fraction of the nanoparticles. Furthermore, a similar study has been done by Etminan-Farooji et al. (2012) in the case of a square cylinder ( $0^{\circ}$ inclined) in nanofluids with two types of nanofluids comprised of $\mathrm{Al}_{2} \mathrm{O}_{3}$ and CuO nanoparticles. Depending upon the type of nanofluids, an enhancement in heat transfer up to $25 \%$ in comparison with pure water can be seen at higher values of the Peclet number and particle volume fraction. From the literature, it is clear that there is no investigation available on the heat transfer phenomena from a square cylinder at different orientation anglesin the uniform flow of nanofluids. This particular shape of the cylinder is important for many process industries (such as novel heat exchangers, various electronic components). Therefore it is important to investigate numerical study of heat transfer from the two-dimensional flow of water base nanofluids over a square cylinder carried out in the laminar steady flow regime. The present numerical results are expressed in terms of isotherm contours (temp, average Nusselt number and discussed the effect of Reynolds number ( $1 \leq \mathrm{Re} \leq 40$ ), nanoparticle volume fraction ( $0 \leq \phi \leq 0.05$ ), nanoparticles ( CuO ), and particle size $d_{n p}$ at constant wall thermal boundary condition (CWT). In the end, the present numerical results have developed the correlation for average Nusselt number for further application.

## 2. PROBLEM STATEMENT

The two-dimensional flow over a square cylinder (size a) to the oncoming uniform stream of nanofluids with velocity $V_{\infty}$ and at temperature, $T_{\infty}$ is shown schematically in Fig. 1. The surface of the square cylinder is maintained at a constant wall temperature $T_{w}$ (CWT) boundary. To make this unconfined flow problem numerically feasible, an artificial cylindrical domain of nanofluids of diameter $D_{\infty}$ is considered in which the square cylinder is placed at the centre of the domain (see Fig. 1b). Over the range of parameters, the flow is assumed to be steady, incompressible, and the laminar flow domain is carried out to minimize the computational efforts. The thermo-physical properties of nanofluids such as density, $\rho_{n f}$, (Pak and Cho, 1998), specific heat, $C_{p, n f}$, (Xuan and Roetzel, 2000), thermal conductivity, $k_{n f}$, (Koo and Kleinstreuer, 2004) and viscosity, $\mu_{n f}$, (Masoumi et al., 2009) are calculated as a function of nanoparticle volume fraction, $\phi$, together with properties of the base fluid and NPs.


Fig. 1. Schematic representation of the problem of flow past and heat transfer from (a) square cylinder, (b) computational domain

Since the thermo-physical properties of the nanofluids are assumed to be temperature-dependent; the velocity and temperature fields are coupled in nature. The viscous dissipation effects are neglected. Within the framework of the above assumptions, the coupled velocity and temperature fields are described by the following mass, momentum, and energy:

- continuity equation:

$$
\begin{equation*}
\Delta V=0 \tag{1}
\end{equation*}
$$

- momentum equation:

$$
\begin{equation*}
\rho_{n f}(V \nabla) V=-\nabla p+\nabla \cdot\left(\mu_{n f} \nabla V\right) \tag{2}
\end{equation*}
$$

- thermal energy equation:

$$
\begin{equation*}
\rho_{n f} C_{n f}(V \nabla) T=\nabla\left(k_{n f} \nabla T\right) \tag{3}
\end{equation*}
$$

The specified boundary conditions are used lower-half of the domain as an inlet condition (i.e., $V_{x}=1$, $V_{y}=0, T=300^{\circ}$ ), no-slip and constant wall temperature on the surface of the square cylinder is defined condition (i.e., $V_{x}=V_{y}=0, T=305^{\circ}$ ), and upper-half of the artificial cylindrical (i.e., $\frac{\partial \varphi}{\partial x}=0$, where $V_{x}$, $V_{y}$ and $T$ ).

Finally, the governing equations subjected to the above boundary conditions are solved in terms of the primitive variables $\left(V_{x}, V_{y}, p, T\right)$ and results are post-processed to evaluate the average value of Nusselt number.

$$
\begin{equation*}
\mathrm{Nu}_{\text {avg }}=\frac{1}{S} \int_{0}^{S} \mathrm{Nud} s \tag{4}
\end{equation*}
$$

The extensive numerical results in the present study are expressed in terms of drag coefficient and Nusselt numbers are expected to be a function of $\operatorname{Re}, \phi, d_{n p}$ and this work endeavours to explore this functional relationship for nanofluids past a heated square cylinder at different orientations.

### 2.1. Numerical solution procedure and computational parameters

The finite element method-based commercial software COMSOL Multiphysics ${ }^{\circledR}$ was used to solve the governing differential equations along with boundary conditions. Detailed numerical descriptions of the solution methodology are available in the literature Kaur et al. (2021). A simulation was deemed to have converged when the relative residuals for the momentum equations were lower than $10^{-7}$. This was also sufficient to ensure that the drag values stabilized at least up to four significant digits. The Lagrangian scheme was used to solve the pressure-velocity coupling. The optimum size of the domain, $D_{\infty} / a=150$ is adequate for this present computational study. The minimum number of a grid with the total number of elements, $N=56298$ was found to be sufficient for the present study over the wide range of conditions considered herein.

## 3. RESULTS AND DISCUSSION

### 3.1. Validation of results

To check the accuracy and confidence of the present numerical results, it is imporment to validate the current numerical methodology. The validation result for forced convection heat transfer phenomena from a square cylinder immersed in water base nanofluids ( CuO ) is represented in Table 1. The average heat transfer coefficient is calculated for the square cylinder in the infinite extent of nanofluids and compared with Etminan-Farooji et al. (2012). The results are seen to be in good agreement, thus, the present numerical results are believed to be reliable to within 3-4\%.

Table 1. Validation of average heat transfer coefficient for a square cylinder in $\mathrm{Al}_{2} \mathrm{O}_{3} /$ water and $\mathrm{CuO} /$ water nanofluids ( $\phi=0.04$ )

| $\mathrm{Pe}=\mathrm{Re} \cdot \mathrm{Pr}$ | $d_{n p}=30 \mathrm{~nm}$ |  | $d_{n p}=100 \mathrm{~nm}$ |  |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Al}_{2} \mathrm{O}_{3}$ | Present | Etminan-Farooji et al. (2012) | Present | Etminan-Farooji et al. (2012) |
| 50 | 2.0131 | 2.1404 | 2.1162 | 2.2275 |
| 100 | 2.6931 | 2.8714 | 2.9847 | 3.0067 |
| 200 | 3.7497 | 3.9593 | 3.9795 | 4.1337 |
| $\mathrm{Pe}=\mathrm{Re} \cdot \mathrm{Pr}$ | $d_{n p}=30 \mathrm{~nm}$ |  | $d_{n p}=100 \mathrm{~nm}$ |  |
| CuO | Present | Etminan-Farooji et al. (2012) | Present | Etminan-Farooji et al. (2012) |
| 50 | 2.1310 | 2.2728 | 2.1086 | 2.2728 |
| 100 | 3.0103 | 3.0610 | 3.0095 | 3.0591 |
| 200 | 4.0899 | 4.1660 | 4.0288 | 4.1660 |

### 3.2. Isotherm contours

Fig. 2 represents the isotherm contours near the square cylinder at $\operatorname{Re}=10$ and $\operatorname{Re}=40$ for two different sizes of NPs, $d_{n p}=30 \mathrm{~nm}$ and $d_{n p}=100 \mathrm{~nm}$ for three different orientation angles $\psi=0,45^{\circ}$ and $75^{\circ}$. At a low value of $\operatorname{Re}(\operatorname{Re}<10)$, the isotherm contours are less crowded near the heated cylinder and weak temperature gradients in the flow domain. Because of the increased amount of inertia force, the crowding of the isotherm contours towards the surface of the cylinder occurs when the Reynolds number ( $\operatorname{Re}>10$ ) increases. This crowding of isotherm contours is seen to be more pronounced on the downstream side of the cylinder due to the large wake formation. The isotherm contours are also crowded with orientation


Fig. 2. Variation of isotherm contour with $\operatorname{Re}, d_{n p}, \phi$ and orientation angle ( $\psi$ )
angle of the cylinder and attained the maximum value at $\psi=45^{\circ}$, after that it is decreasing to $\psi=90^{\circ}$ (not indicated in Fig. 1). Crowding of isotherm contours indicated that reaction that occurs near the cylinder suppresses the isotherms close to the surface of the cylinder thereby resulting in the thin thermal boundary layers. Next, with nanofluids, such as the addition of $N P s$ in the base fluids, the thermal boundary layers are seen to be much thinner as compared to the pure fluids (water), and as a result, the isotherm contours are much denser near the surface of cylinder irrespective of the values of $\operatorname{Re}, \psi$ and $d_{n p}$. This effect becomes more effective as the value of the Reynolds number progressive increases. Hence, one can expect some enhancement in heat transfer in the presence of nanofluids as compared to pure fluids. Indeed this assertion is borne out by the Nusselt number results presented in the next section.

### 3.3. Effect of orientation of cylinder

Figure 3 represents the orientation effect of square cylinder on average Nusselt Number, $\mathrm{Nu}_{\text {avg }}$, at different volume fractions ( $\phi=0.1$ and $\phi=0.05$ ), sizes of NPs ( $d_{n p}=30 \mathrm{~nm}$ and $d_{n p}=100 \mathrm{~nm}$ ) and two values of Reynolds number $(\operatorname{Re}=10$ and $\operatorname{Re}=40)$. At the low value of $\operatorname{Re}(\operatorname{Re}=10)$, the orientation effect of the cylinder $\mathrm{Nu}_{\text {avg }}$ is almost indistinguishable due to the low inertial effect, and consequently, lower heat transfer occurred for all orientations of the cylinder. As the value of $\operatorname{Re}$ increases $(\operatorname{Re}=40)$, the orientation effect on the value of $\mathrm{Nu}_{\text {avg }}$ is visible due to the higher inertial effect which directly influences the thermal boundary layer irrespective of NPs size. It can be seen that the thermal boundary layer for the square cylinder at $45^{\circ}$ orientation is smaller especially near the upstream and downstream side of the cylinder which is directly related to $\mathrm{Nu}_{\mathrm{avg}} \sim 1 / \delta_{t h}$. This is also consistent with the isotherm contour shown in Fig. 2. It is seen that the heat transfer from the square cylinder at $45^{\circ}$ orientation generated large heat transfer as compared to other orientations of the cylinder.


Fig. 3. Variation of average Nusselt number with orientation angle at different values of $\operatorname{Re}, \phi$ and $d_{n p}$

Next, the effect of $\phi$ and $d_{n p}$ can be seen for $45^{\circ}$ orientated square cylinders. It is better to represent the variation of heat transfer enhancement ( $E=\mathrm{Nu}_{\text {avg }} / \mathrm{Nu}_{a v g, \phi=0}$ ) with NPs volume fraction, $\phi$, is represented at the different sizes of nanoparticles, $d_{n p}$. From Fig. 4, it is observed that at the small size of NPs (such as $30 \mathrm{~nm}, 50 \mathrm{~nm}$ ), the value of $E$ increases with $\phi$ up to the optimum value, after that it starts to deteriorate. As the size of NPs are large (such as 60 nm and 100 nm ), the value of $E$ gradually increases with $\phi$ at a fixed value of Re. This is due to the effective Reynolds number being higher than its global value of Reynolds number in the case of nanofluids containing large NPs of 60 nm in size, thereby rate heat transfer increases with $\phi$.


Fig. 4. Variation of $E$ (for 45 inclined square cylinder) with $\phi$ of $\mathrm{CuO} N P s$

From the application point of view, the average Nusselt number for $45^{\circ}$ orientated square cylinder (optimum orientation of the cylinder) is correlated by using a simple expression for different sizes of NPs ( 30 nm to 100 nm ) which is used to find out the intermediate values of $\phi$ and $\operatorname{Re}$ over the ranges of condition $1 \leq \operatorname{Re} \leq 40$ and $0.00 \leq \phi \leq 0.05$.

$$
\begin{equation*}
\mathrm{Nu}_{\mathrm{avg}}=\frac{a}{\operatorname{Re}} \frac{\left(1+b \mathrm{Re}^{c}\right)}{\left(1+e \phi^{g}\right)}(1+d \phi) \tag{5}
\end{equation*}
$$

In Eq. (5), the values of the fitting parameters obtained through the non-linear regression analysis are given in Table 2.

Table 2. The fitting parameter value for the average Nusselt number

| $N P s$ | $d_{n p}$ | $a$ | $b$ | $c$ | $d$ | $e$ | $g$ | Avg \% <br> error | Max \% <br> error |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| CuO | 30 nm | 26.981 | 0.765 | 0.734 | -15.279 | -9.2171 | 0.771 | 8.53 | 23.64 |
|  | 100 nm | 20.925 | 1.056 | 0.793 | -4.7873 | -5.7873 | 0.913 | 6.87 | 21.73 |

## 4. CONCLUSIONS

A numerical study of laminar, incompressible $2 D$ heat transfers of CuO -water-based nanofluids past a heated square cylinder at different orientations $\left(0-90^{\circ}\right)$ in an infinite domain of nanofluids was conducted. For Reynolds number (1-40), nanoparticle volume fraction ( $0-0.05$ ), size of the NPs ( $30-100 \mathrm{mn}$ ), the orientation of square cylinder, the following conclusions can be drawn.

- The crowding of isotherm contoursnear the cylinder increased with the orientation of the square cylinder and attained the maximum crowding at $45^{\circ}$ orientated square cylinder which led to the decrease of the thermal boundary layer.
- The average Nusselt number of a square cylinder at low Reynolds number is indistinguishable regardless of the values of $\phi$ and $d_{n p}$, whereas, in the case of higher Reynolds number, there is an optimum value of at $45^{\circ}$ orientated square cylinder. It starts to deteriorate with the further increasing the orientation angle.
- Finally developed the simple correlation for the average Nusselt number is $\mathrm{Nu}_{\text {avg }}$ in the given range of conditions.


## SYMBOLS

| a | side length, m |
| :--- | :--- |
| $d_{n p}$ | diameter of nanoparticles, nm |
| $D_{\infty}$ | diameter of nanoparticles, nm |
| $h$ | heat transfer coefficient, dimensionless |
| $k_{n f}$ | thermal conductivity of nanofluid, $\mathrm{W} \cdot \mathrm{m}^{-1} \cdot \mathrm{~K}^{-1}$ |
| $k_{b f}$ | thermal conductivity of base fluid, $\mathrm{W} \cdot \mathrm{m}^{-1} \cdot \mathrm{~K}^{-1}$ |
| $\mathrm{Nu}_{\text {avg }}$ | average Nusselt number, dimensionless |
| $N$ | total number of elements in the computational domain, dimensionless |
| Re | Reynolds number, dimensionless |
| $T_{\infty}$ | temperature at inlet, K |
| $T_{w}$ | temperature at wall, K |
| $V_{\infty}$ | velocity at inlet, $\mathrm{m} / \mathrm{s}$ |
| $\Psi$ | orientation angle of $\mathrm{squrae} \mathrm{cylinder} deg$, |
| $\rho$ | density of fluid, $\mathrm{kg} \cdot \mathrm{m}^{-3}$ |
| $\mu$ | viscosity of fluid, $\mathrm{Pa} \cdot \mathrm{s}$ |
| $\rho_{n p}$ | density of nanopartcile, $\mathrm{kg} \cdot \mathrm{m}^{-3}$ |
| $\rho_{b f}$ | density of base fluid, $\mathrm{kg} \cdot \mathrm{m}^{-3}$ |
| $\rho_{n f}$ | density of nano fluid, $\mathrm{kg} \cdot \mathrm{m}^{-3}$ |
| $\mu_{n f}$ | viscosity of nano fluid, $\mathrm{Pa} \cdot \mathrm{s}$ |
| $\mu_{b f}$ | viscosity of base fluid, $\mathrm{Pa} \cdot \mathrm{s}$ |
| $n f$ | nano fluid |
| $b f$ | base fluid |

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